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Spark Plasma Sintered NdFeB-based Nanocomposite Hard Magnets with Enhanced Magnetic Properties

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1. Introduction

Since the discovery of 2:14:1 permanent magnets (PMs) in the early 80’s, the nanocrystalline and nanocomposite RE-Fe-B PMs attracted extensive attention both from the academia and industry because of their enhanced magnetic performances and potential applications (Croat et al., 1984; Givord et al., 1984; Hadjipanayis et al., 1984; Onodera et al., 1984; Yamauchi et al., 1985; Jha & Davies, 1989; Pinkerton, 1991; Hadjipanayis, 1999; Liu et al., 2008; Fukagawa et al., 2010). Whereas the enhanced remanence and energy product of nanocrystalline Nd-Fe-B PMs result from the exchange coupling between magnetically hard grains (Manaf et al., 1993), for the nanocomposite PMs the superior magnetic properties are the result of the fine mixture of RE$_2$Fe$_{14}$B hard magnetic and Fe-based soft magnetic grains, which are exchange coupled (Kneller & Hawig, 1991). Typically, the exchange coupling in nanophase magnets tends to reduce the coercivity (Mendoza-Suarez et al., 2000) and the energy product decreases with the volume fraction of the magnetic phase, which are well-known challenges in the processing of PMs. The recent work proved the important role played by magnetostatic interactions in the increase of the nucleation and coercive field values of nanocomposite permanent magnets in the detriment of the exchange coupled interactions (Gabay et al., 2006; Marinescu et al., 2008).

The properties of nanocomposite PMs are strongly influenced by a number of process parameters like the composition, preparation method, annealing conditions, distribution of soft and hard magnetic nanograins (Cui et al., 2005). The large values of the remanence are related to the strength of the exchange interactions between the soft and hard magnetic grains, thus the reduction of the grains below the size of the hard magnetic phase domain walls is essential, but they should not decrease below a critical value, dependent on the nanocomposite composition, which causes the reduction of the magnetocrystalline anisotropy of the hard magnetic phase and, consequently, the drastic decrease of the coercive field. Nanocomposite magnets are expected to have maximum energy products as high as 120 MGOe, when the soft and hard magnetic phases are arranged in the proper way and the exchange interactions optimized (Skomski & Coey, 1993). However, the theoretical value of the maximum energy product of NdFeB magnets is calculated to be 512 kJ/m$^3$ (64 MGOe) (Sagawa et al., 1985).

There are two principal manufacturing routes to prepare RE-TM-B nanocomposite PMs: (i) the classical powder metallurgy or sintered magnet process (Sagawa et al., 1984; Durst &
Kronmüller, 1987; Kaneko et al., 2006) and (ii) the rapid solidification or bonded magnet process (McGuiness et al., 1992a; McGuiness et al., 1992b; Tomka et al., 1995; Folks et al., 1995; Zhang & Xiong, 2009). Sintered magnets are prepared from magnetically aligned powders into dense blocks, which are then heat treated, cut to shape, surface treated and magnetized. Bonded magnets are prepared from milled melt-spun ribbons mixed with a polymer or a resin and either compressed or injection moulded. The bonded magnets offer less flux than sintered ones but can be shaped into intricately shaped parts and their eddy current losses are significantly smaller. The melt-spun ribbons can be also hot pressed into fully dense isotropic magnets, and then forged or extruded into high energy anisotropic magnets.

About a decade ago a new technique appeared as an alternative to the conventional sintering methods. Spark Plasma Sintering (SPS) method can consolidate Nd-Fe-B powders at relatively lower temperatures in a short period of time, and the achieved densities of compacted materials can go up to 99% (Liu et al., 1999; Saito, 2001; Ono et al., 2003). But the main and the most important advantage of SPS technique is the high sintering speed, which can effectively hinder the grains growth over the critical nanosize. However, the structure and properties of the compacted materials are dependent on several process parameters, such as the sintering temperature and the applied pressure.

In this following, we will present our results on the preparation and magnetic properties of a few types of Spark Plasma Sintered NdFeB-based nanocomposite magnets prepared from mixtures of ball-milled powders of stoichiometric and non-stoichiometric Nd-Fe-B nanocrystalline melt-spun ribbons and Fe-based amorphous and crystalline materials. The novelty, compared with the already published work on SPS nanocomposite magnets, consists of the use for the first time ever of Fe and Fe-Co-based amorphous/nanocrystalline materials as soft magnetic components.

2. Nd-Fe-B melt spun ribbons

Stoichiometric Nd_{12}Fe_{82}B_{6} (Nd12), sub-stoichiometric Nd_{8}Fe_{86}B_{6} (Nd8) and Nd_{10}Fe_{84}B_{6} (Nd10), and Nd-rich Nd_{15}Fe_{79}B_{6} (Nd15) and Nd_{16}Fe_{78}B_{6} (Nd16) master alloys have been prepared by arc melting using a mixture of pure elements in argon atmosphere. The alloys have been re-melted a few times for homogenization and used further to prepare amorphous and nanocrystalline ribbons having thicknesses of 25 µm and widths of 3-5 mm by melt-spinning technique, as shown in Figs. 1 and 2. Whereas the Nd8 XRD pattern (Fig. 1) indicates only the presence of the Fe_{3}B peak superposed on the broad one of the amorphous phase, the Nd10 ribbon consists of a mixture of αFe, Fe_{3}B and amorphous phase. The stoichiometric Nd12 melt-spin ribbons are nanocrystalline in the as-quenched state, but the size of the grains as determined by using the Scherrer formula is around 100 nm, and the presence of the residual amorphous matrix is still present. The DSC curves shown in Fig. 2 indicate also the decrease of the amount of amorphous phase with the increase of the Nd content, and for samples with Nd content over 12 at.% the amorphous phase contribution can not be observed anymore. Whilst the Nd-rich ribbons (Nd_{15}Fe_{79}B_{6} and Nd_{16}Fe_{78}B_{6}) have been further used in the as-quenched state as precursors to prepare powders by high-energy ball-milling, the stoichiometric Nd_{12}Fe_{82}B_{6} and sub-stoichiometric Nd_{8}Fe_{86}B_{6} and Nd_{10}Fe_{84}B_{6} ones have been first annealed to achieve the optimum nanocrystalline structure and then ball-milled.
Spark Plasma Sintered NdFeB-based Nanocomposite Hard Magnets with Enhanced Magnetic Properties

Fig. 1. XRD patterns of Nd-Fe-B melt-spun ribbons in the as-quenched state as a function of Nd content.

Fig. 2. DSC curves for Nd-Fe-B melt-spun ribbons vs. Nd content.

The variation of the magnetic characteristics of Nd8, Nd10 and Nd12 melt-spun ribbons as a function of annealing temperature and time are shown in Fig. 3 and Table 1. The coercive field increases with the increase of the Nd content as the result of the increase of the amount of 2:14:1 hard magnetic phase, which in the same time causes a slight decrease of the saturation magnetization and remanence. The coercive field of Nd8 ribbons starts increasing with the annealing temperature, when the 2:14:1 and Fe₃B phases are precipitating. The further increase of the annealing temperature keeps the coercive field almost constant in the temperature range 610-650°C and then starts decreasing slowly due to the increase of the soft magnetic grains. For Nd10 melt-spun ribbons, the presence of αFe and Fe₃B phases in the as-quenched state results only in the very slight variation of the magnetic characteristics, whereas for the already nanocrystalline Nd12 melt-spun ribbons in the as-quenched state the annealing is refining the nanograins structure and consequently the coercive field increases whereas the saturation magnetization and remanence remain almost unchanged. From Fig. 3 one can conclude that the optimum annealing temperature is 630°C for Nd8 and Nd10 melt-spun ribbons, and 610°C for Nd12 melt-spun ribbons, respectively.
Fig. 3. The variation of the coercive field and remanence as a function of the annealing temperature ($T_a$) for Nd8, Nd10 and Nd12 melt-spun ribbons.

Keeping the annealing temperature constant we varied the annealing time and the obtained results are given in Table 1.

<table>
<thead>
<tr>
<th>Composition</th>
<th>Annealing time (min.)</th>
<th>$H_c$ (kOe)</th>
<th>$4\pi M_s$ (kG)</th>
<th>$4\pi M_r$ (kG)</th>
<th>$M_r/M_s$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Nd$<em>8$Fe$</em>{86}$B$_6$ ($T_a=630^\circ$ C)</td>
<td>10</td>
<td>3.6</td>
<td>13.2</td>
<td>7.6</td>
<td>0.57</td>
</tr>
<tr>
<td></td>
<td>20</td>
<td>4.2</td>
<td>13.1</td>
<td>8.4</td>
<td>0.63</td>
</tr>
<tr>
<td></td>
<td>30</td>
<td>2.8</td>
<td>13.1</td>
<td>7.8</td>
<td>0.60</td>
</tr>
<tr>
<td>Nd$<em>{10}$Fe$</em>{84}$B$_6$ ($T_a=630^\circ$ C)</td>
<td>10</td>
<td>5.0</td>
<td>11.3</td>
<td>8.6</td>
<td>0.75</td>
</tr>
<tr>
<td></td>
<td>20</td>
<td>5.1</td>
<td>11.7</td>
<td>8.4</td>
<td>0.72</td>
</tr>
<tr>
<td></td>
<td>30</td>
<td>5.2</td>
<td>11.0</td>
<td>8.1</td>
<td>0.73</td>
</tr>
<tr>
<td>Nd$<em>{12}$Fe$</em>{82}$B$_6$ ($T_a=610^\circ$ C)</td>
<td>10</td>
<td>5.0</td>
<td>10.8</td>
<td>8.0</td>
<td>0.73</td>
</tr>
<tr>
<td></td>
<td>20</td>
<td>6.2</td>
<td>11.8</td>
<td>8.5</td>
<td>0.74</td>
</tr>
<tr>
<td></td>
<td>30</td>
<td>5.0</td>
<td>10.6</td>
<td>7.6</td>
<td>0.72</td>
</tr>
</tbody>
</table>

Table 1. Magnetic characteristics of Nd-Fe-B melt-spun ribbons vs. annealing time.

The increase of the annealing time from 10 to 20 min. for Nd8 melt-spun ribbons enhances the magnetic characteristics because of the achievement of the full nanocrystalline structure. The further increase of the $T_a$ to 30 min. produces the decoupling of the soft and hard magnetic nanograins due to the increase over the critical size of the soft magnetic nanograins. Due to the already existing nanocrystalline structure in the as-quenched state,
the annealing of the Nd12 melt-spun ribbons at 610°C for 10 min. is enough to refine the nanograins structure and to increase the number of exchange interactions. By increasing even further the annealing time the nanograins are growing excessively and the magnetic properties are altered. The magnetic characteristics of the Nd10 melt-spun ribbons are not changing significantly with the annealing time modification, most probably because of the precipitation in the as-quenched state of soft magnetic \( \alpha \)Fe and Fe\(_3\)B phases.

3. Nd-Fe-B powders

Nd-Fe-B melt-spun ribbons have been cut in small pieces and then ball-milled in 2 stainless steel vials with 15 stainless steel balls in each vial. It is well known that the characteristics of the sintered magnets are strongly dependent on the magnetic characteristics of the ball-milled powders. The resulting Nd-Fe-B powders have sizes ranging from 20 to 100 μm. We’ve noticed that by increasing the milling time the magnetic properties of the powders are deteriorating. This behaviour could have 2 distinct causes: (1) easier oxidation of the powders when they are becoming smaller for longer milling times (despite the fact that milling is done in protective atmosphere), and (2) the precipitation of more Fe-based soft magnetic phase(s) in the detriment of 2:14:1 hard magnetic phase. To minimize the effect of the magnetic characteristics alteration the melt-spun ribbons were milled only for 4 min. Nd-Fe-B powders, and especially those with larger contents of Nd (≥12 at.%), are very brittle, so smaller milling times can provide the powders in the desired range of sizes. In Table 2 are presented comparatively the magnetic characteristics for stoichiometric and sub-stoichiometric Nd-Fe-B melt-spun ribbons and the resulting ball-milled powders with sizes ranging from 20 to 100 μm.

<table>
<thead>
<tr>
<th>Composition</th>
<th>Sample type</th>
<th>( iH_c ) (kG)</th>
<th>( 4\pi M_s ) (kG)</th>
<th>( 4\pi M_r ) (kG)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Nd(<em>8)Fe(</em>{86})B(_6)</td>
<td>Melt-spun ribbon</td>
<td>4.2</td>
<td>13.2</td>
<td>8.4</td>
</tr>
<tr>
<td></td>
<td>Powders</td>
<td>4.3</td>
<td>12.0</td>
<td>7.3</td>
</tr>
<tr>
<td>Nd(<em>{10})Fe(</em>{84})B</td>
<td>Melt-spun ribbon</td>
<td>5.0</td>
<td>11.9</td>
<td>8.9</td>
</tr>
<tr>
<td></td>
<td>Powders</td>
<td>5.2</td>
<td>11.0</td>
<td>7.5</td>
</tr>
<tr>
<td>Nd(<em>{12})Fe(</em>{82})B(_6)</td>
<td>Melt-spun ribbon</td>
<td>6.3</td>
<td>11.2</td>
<td>8.4</td>
</tr>
<tr>
<td></td>
<td>Powders</td>
<td>7.7</td>
<td>10.3</td>
<td>6.6</td>
</tr>
</tbody>
</table>

Table 2. Magnetic characteristics for annealed melt-spun ribbons and the resulting powders.

One can notice the increase of the coercive field and the slight decrease of the saturation magnetization and remanence for the resulting powders. The increase of the coercive field is more pronounced for the powders with larger Nd contents, most probably as a result of the decoupling between the soft and hard magnetic phases, as shown in Fig. 4. For Nd8 and Nd10 powders the exchange coupling is improved (the shoulder in the second quadrant disappears) compared with the precursor melt-spun ribbons, confirming the previously observed results for milled powders used to make bonded magnets (McGuiness et al., 1992a; McGuiness et al., 1992b; Tomka et al., 1995; Folks et al., 1995; Zhang & Xiong, 2009). A consequence of better exchange coupling between the soft and hard magnetic grains is also the small reduction of the remanence for Nd8 and Nd10 ball-milled powders. It is important to note also that the magnetization hysteresis loop for Nd12 powders is not saturated, but clearly shows the presence of 3 magnetic phases: soft magnetic \( \alpha \)Fe and Fe\(_3\)B and hard magnetic 2:14:1 phase, whereas the melt-spun ribbons consist only of \( \alpha \)Fe soft magnetic and 2:14:1 hard magnetic phase.
Fig. 4. Magnetic hysteresis curves for Nd10 and Nd12 annealed melt-spun ribbons (at 630°C for 20 min. for Nd10 and 610°C for 10 min. for Nd12), respectively, and the resulting ball-milled powders with sizes ranging from 20 to 100 μm.

Fig. 5 presents the hysteresis loops for Nd15Fe79B6 melt-spun and ball-milled powders of 63-80 μm as well as for Nd16Fe78B6 melt-spun and powders in the range size of 80-100 μm.

As one can see the hysteresis curves of powders are constricted compared with the melt-spun ribbon’s ones, which might be an indication that during ball-milling the precipitation of more Fe-based phases is favoured in the detriment of 2:14:1 phase and the decoupling is more evident. Additionally, the magnetic behaviour of Nd-rich alloys powders is strongly dependent on the powders size compared with the stoichiometric and sub-stoichiometric compositions for which we didn’t notice such a behaviour. The optimum magnetic properties are achieved for Nd15 powders in the 63-80 μm range and respectively for Nd16 powders with sizes between 80 and 100 μm, as shown in Fig. 6. Most probably, this behaviour is a collective process which involves both the surface oxidation effects (despite the fact that the ball-milling is done in protective Ar atmosphere) and the grains size refinement. It is worthwhile to note the high values of the coercive field obtained for both Nd15 and Nd 16 melt-spun nanocrystalline ribbons and ball-milled powders of about or over 20 kOe, but also the moderate values of the remanence which are a clear proof that the
predominant phase is 2:14:1, whereas the volume of Fe-based soft magnetic phases is much smaller compared with Nd8, Nd10 and Nd12 samples. One should also notice that the applied field of 3 T used in this work is not enough to saturate either the melt-spun ribbons or ball-milled powders.

Fig. 6. Demagnetization curves for Nd15 and Nd16 ball-milled powders in the size range 10-150 μm.

4. Soft magnetic materials used as precursors for SPS nanocomposite PMs

For this study we have chosen a number of soft magnetic materials to be used as precursor for preparing SPS nanocomposite magnets: (1) Fe_{77.5}Si_{7.5}B_{15} (FeSiB) \( (4πM_s \sim 16 \text{ kG}) \) and (2) Co_{68.25}Fe_{4.5}Si_{12.25}B_{15} (CoFeSiB) \( (4πM_s \sim 8 \text{ kG}) \) conventional amorphous wires (CAW) of 100-125 μm in diameter prepared by in-rotating water quenching method (Ohnaka, 1985; Vazquez & Chen, 1995; Chiriac et al., 1998); (2) Co_{68.2}Fe_{4.5}Mn_{0.1}Si_{12.25}B_{15} (CoFeMnSiB) amorphous melt-spun ribbons \( (4πM_s \sim 8 \text{ kG}) \) (Hilzinger & Kunz, 1980; Makino et al., 1990; Chiriac et al., 1997); (3) Fe_{73.5}Nb_{3}Cu_{13.5}B_{9} (FINEMET) nanocrystalline melt-spin ribbons \( (4πM_s \sim 12.4 \text{ kG}) \) (Yoshizawa et al., 1988); (4) Fe_{2}Co crystalline melt-spin ribbons \( (4πM_s \sim 22 \text{ kG}) \); and (5) commercial Fe micropowders \( (4πM_s \sim 20 \text{ kG}) \) of about 10 μm. The studies done up to know on Nd-F-B nanocomposite bulk magnets (sintered or bonded) used either Fe (Wei et al., 1994; McCormick et al., 1998; Lewis et al., 2001; Rada et al., 2005; Niu et al., 2007; Bhame et al., 2010; Cha et al., 2010) or Fe_{2}Co micro or nanopowders (Wei et al., 1994) to increase the remanence and saturation magnetization of the respective PMs, whereas Fe, Fe-Co and Fe-B have been used in combination with Nd-Fe-B to prepare multilayered structures by deposition methods (Shindo et al., 1996; Cui & O’Shea, 2004; Ao et al., 2005; Kato et al., 2005). Additionally, other attempts to prepare nanocomposite PMs used combinations of Nd-Fe-B and Sm-Co-based particles (Talijan, 2006; Betancourt & Davies, 2010). However, excepting the multilayered structures obtained by deposition, most of the other studies are related to nanocomposite magnets prepared by milling of the starting Nd-Fe-B alloys followed by mixing with Fe-based soft magnetic materials and then annealing of the full composite, and not by milling the melt-spin ribbons in the as-quenched state or after annealing. Our goal is to use soft magnetic high saturation magnetization materials in combination with Nd-Fe-B ball milled powders.
FeSiB and CoFeSiB CAW have been cut in small pieces and then ball-milled in argon atmosphere for 6 h. The resulting micropowders have sizes between 5 and 150 µm, and their magnetic properties are almost independent of the powders size (Fig. 7) compared with Nd-Fe-B ball milled powders. The amorphous state is preserved even after milling for 6 h, as shown by the structural studies.

![Fig. 7. Magnetic hysteresis curves for FeSiB and CoFeSiB as-quenched conventional amorphous wires (CAW) and the resulting ball-milled powders.](image)

The same behaviour was observed for amorphous CoFeMnSiB and FINEMET powders prepared from amorphous and respectively nanocrystalline melt-spun ribbon precursors. The amorphous FINEMET ribbons have been annealed at 560°C for 1 h to achieve the optimum αFe nanograins size (~ 10 nm) and then ball-milled for 20 h. The same milling time was used to prepare powders from CoFeMnSiB amorphous ribbons. After milling the nanocrystalline structure of FINEMET ribbons and respectively the amorphous structure of CoFeMnSiB melt-spun ribbons are preserved, and the ball-milled powders show almost identical XRD patterns.

X-ray diffraction pattern of Fe₂Co (FeCo) melt-spun ribbons of 25 µm thick prepared by melt-spinning reveals only the bcc lines. The same structure is preserved after ball-milling the ribbons for 24 h. The magnetic characteristics measured for the precursor melt-spun ribbons and ball-milled powders are almost similar, with saturation magnetization of about 22 kG.

The commercial Fe micropowders of 10 µm used in this study to prepare SPS nanocomposite magnets have a saturation magnetization of ~20 kGs, a little bit lower than the saturation magnetization of bulk Fe (~ 22 kG), indicating the presence of some traces of Fe oxides on their surface as confirmed by XRD measurements (very small peaks corresponding to Fe₂O₃ can be identified along with the bcc lines of Fe).

5. Spark Plasma Sintered (SPS) nanocomposite magnets

Nd-Fe-B and soft magnetic materials powders have been mixed in different ratios (80:20; 90:10; 92:8; 94:6; 95:5 and 96:4, respectively), placed into a carbon die, and then sintered under vacuum by SPS. The temperature of the specimen was increased from room
temperature to the consolidating temperature of 600ºC for 300 s and then held at the consolidating temperature for 300 s by applying a pulsed electric current of 800 A at intervals of 2.4 ms. The temperature was measured using a thermocouple placed on the surface of the sample. The specimens have been prepared at different pressures ranging from 50 to 80 MPa, applied during the heating. After the SPS consolidation, the samples have been cooled down to room temperature without applying pressure. The specimens were 10 mm in diameter and 2-5 mm in thickness. Their structural and magnetic properties have been investigated afterwards by means of X-ray diffractometry (BRUKER AXS D8-ADVANCE XRD), scanning electron microscopy combined with EDS (Carl Zeiss CrossBeam Neon 40 EsB), and vibrating sample magnetometry (LakeShore 7410 VSM).

Fig. 8 shows the XRD patterns for a number of Nd-Fe-B/soft magnetic material SPS nanocomposite magnets. One can notice that after spark plasma sintering the 2:14:1 phase is preserved and additional peaks corresponding to $\alpha$Fe(Co) and Fe$_3$B soft magnetic phases can be easily identified. However, the amount of soft magnetic phases is larger compared with the Nd-Fe-B melt-spun precursors due to the supplementary addition of Fe and Co-Fe based soft magnetic materials. As smaller the amount of Nd in the precursor Nd-Fe-B powders, as larger the amount of 2:14:1 phase is, contrary to the expectations. This means the addition of soft magnetic powders is contributing to the precipitation of a larger amount of 2:14:1 phase in the nanocomposite samples.

![XRD patterns](https://www.intechopen.com)

**Fig. 8.** XRD patterns of Nd-Fe-B/soft magnetic powders SPS nanocomposite magnets.

The formation of a larger amount of 2:14:1 hard magnetic phase is also demonstrated by the thermomagnetic curves presented in Fig. 9. It is worth noticing that the largest saturation magnetization is achieved for Nd12-4FeCo SPS nanocomposites and also a significant increase is observed for all Nd-Fe-B/FeCo nanocomposite samples as expected. The presence of $\alpha$Fe is observed for all nanocomposite samples ($T_C \sim 770^\circ$C) and the largest amount of soft magnetic $\alpha$Fe is present in Nd-Fe-B/Fe magnets.
Fig. 9. M-T curves for Nd-Fe-B/soft magnetic powders SPS nanocomposite magnets. The thermomagnetic curves for Nd15 and Nd16 melt-spun ribbons are shown for comparison. The values of the magnetic characteristics for several Nd-Fe-B/FeSiB and Nd-Fe-B/CoFeSiB nanocomposite magnets prepared by SPS are shown in Table 3.

<table>
<thead>
<tr>
<th>Sample</th>
<th>$4\pi M_r$ (kG)</th>
<th>$H_c$ (kOe)</th>
<th>$(BH)_{max}$ (MGOe)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Nd15Fe76B6</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Melt-spun ribbons</td>
<td>5.0</td>
<td>20.30</td>
<td>6.2 (25.4)</td>
</tr>
<tr>
<td>Ball-milled powders (63-80 µm)</td>
<td>6.3</td>
<td>19.01</td>
<td>9.9 (29.9)</td>
</tr>
<tr>
<td>NdFeB-4FeSiB / 70 MPa</td>
<td>5.8</td>
<td>16.02</td>
<td>8.4 (23.2)</td>
</tr>
<tr>
<td>NdFeB-5FeSiB / 50 MPa</td>
<td>6.5</td>
<td>17.85</td>
<td>10.6 (29.0)</td>
</tr>
<tr>
<td>NdFeB-10FeSiB / 50 MPa</td>
<td>5.9</td>
<td>16.97</td>
<td>8.7 (25.0)</td>
</tr>
<tr>
<td>NdFeB-5CoFeSiB / 50 MPa</td>
<td>6.2</td>
<td>17.45</td>
<td>9.6 (27.0)</td>
</tr>
<tr>
<td>NdFeB-10CoFeSiB / 50 MPa</td>
<td>5.4</td>
<td>15.13</td>
<td>7.3 (20.4)</td>
</tr>
</tbody>
</table>

| Nd16Fe78B6 |
| Melt-spun ribbons | 4.82 | 20.44 | 5.8 (24.6) |
| Ball-milled powders (80-100 µm) | 5.70 | 16.17 | 8.1 (23.0) |
| NdFeB-4FeSiB / 50 MPa | 6.27 | 17.41 | 9.9 (27.3) |
| NdFeB-5FeSiB / 50 MPa | 5.87 | 16.97 | 8.7 (24.9) |
| NdFeB-4CoFeSiB / 50 MPa | 6.11 | 17.41 | 9.3 (26.6) |
| NdFeB-5CoFeSiB / 50 MPa | 5.99 | 17.02 | 9.0 (25.5) |

Table 3. The magnetic characteristics of several Nd-Fe-B/FeSiB and Nd-Fe-B/CoFeSiB nanocomposite magnets prepared by SPS method, as a function of the hard to soft materials ratio and the applied pressure during consolidation. The $(BH)_{max}$ values in the brackets are calculated for a packing degree of 98% (Lupu et al., 2009) and for an applied field of 50 kOe. The largest coercive field is obtained for melt-spun ribbons, whilst the powders show larger values of the remanence. By consolidation through the SPS technique the maximum energy product generally increases for optimal processing parameters, the highest increase of about 10% being obtained for Nd16-4FeSiB nanocomposite magnet consolidated at 50 MPa. The degree of compaction for the below mentioned samples is reaching 65%, and if we extrapolate the data to an apparent density of 98% (Saito et al., 2004) and we assume the
samples are saturated in an external field of 50 kOe then the maximum energy product for our isotropic SPS nanocomposite magnets is reaching values of 25-29 MGOe, comparable with those obtained up to now for anisotropic nanocomposite magnets (Saito, 2001; Yue at al., 2003).

Fig. 10(a) presents comparatively the $dM/dH=f(H)$ curves for Nd15 ball-milled powders and 2 SPS nanocomposite magnets (Nd15-5FeSiB and Nd15-5CoFeSiB), both consolidated at 50 MPa, whilst Fig. 10(b) shows the variation of $dM/dH$ with the soft magnetic material content for the same consolidating pressure.

![Fig. 10. dM/dH=f(H) curves for Nd15 ball-milled powders and SPS consolidated nanocomposite magnets as a function of added soft magnetic material (Lupu et al., 2009).](image)

One can notice that despite the fact that the powders exhibit larger coercive field, the coupling between the soft and hard magnetic phases is improved for nanocomposite magnets. The strength of the decoupling is indicated by the second shoulder of $dM/dH=f(H)$ curves located nearby Y axis, which is very small for SPS consolidated nanocomposite magnets, but also by the height of the first peak. In this case, it is hard to believe that only the intergranular exchange coupling is responsible for such behaviour, but most probably a collective process of both exchange and magnetostatic interactions between the powders of the 2 materials forming the nanocomposite (Gabay et al., 2006).

The same behavior is observed for Nd16 powders and nanocomposite magnets, as indicated in Fig. 11. As expected, the exchange coupling plays the predominant role in melt-spun nanocrystalline ribbons and the presence of the second shoulder is missing, whereas for the composites the magnetostatic interactions becomes more predominant (see also Fig. 12). However, the percentage of exchange interactions contribution to the macroscopic response of the nanocomposite systems is still very important, considering that usually the magnetostatic interactions are reducing the coercive field values.

The maximum energy product increases with about 13% for Nd15/soft magnetic material and Nd16/soft magnetic material nanocomposite SPS magnets. The optimum sintering temperature is 50 MPa.

The interactions between the nanograins of SPS nanocomposite magnets were examined by using the $\delta M$ plot, defined as:

$$\delta M(H) = m_d(H) - [1 - 2m_r(H)],$$
where \( m_d \) is the reduced demagnetization remanence and \( m_r \) is the reduced isothermal remanence. According to Wohlfarth’s model (Wohlfarth, 1958), a positive \( \delta M \) peak indicates an exchange coupling process, while a negative \( \delta M \) peak indicates magnetostatic interactions between the magnetically hard and soft phases. A sharp peak indicates strong exchange interactions between the constituent grains, and a less pronounced peak reveals weak exchange intergranular interactions.

Fig. 11. \( dM/dH=f(H) \) curves for Nd16 melt-spun ribbons, ball-milled powders and SPS consolidated nanocomposite magnets as a function of added soft magnetic material (Lupu et al., 2009).

Fig. 12. Magnetic hysteresis loops for Nd-Fe-B nanocomposite SPS magnets. The M-H loops for Nd15 and Nd16 melt-spun ribbons and precursor ball-milled powders are also plotted for comparison.

In Fig. 13 are presented the \( \delta M \) curves for Nd15/FeSiB and Nd15/CoFeSiB SPS nanocomposites with different contents of soft magnetic material. The magnets with less soft magnetic component show larger \( \delta M \) values, i.e. stronger exchange interactions between the nanograins, confirming the large values of the remanence and maximum energy product presented in Table 3.
Spark Plasma Sintered NdFeB-based Nanocomposite Hard Magnets with Enhanced Magnetic Properties

Fig. 13. $\delta M$ curves for Nd$_{15}$Fe$_{79}$B$_6$/FeSiB and Nd$_{15}$Fe$_{79}$B$_6$/CoFeSiB SPS nanocomposite magnets as a function of the soft magnetic component content. Powders with sizes ranging from 63 to 80 $\mu$m were used to prepare the magnets and a pressure of 50 MPa was applied during SPS process.

The sintering pressure not only displaces the powders to fill the empty spaces but introduces also some frictions between particles. The generated heat during the filling of the empty volume as well as the heat generated through Joule effect when the electrical current is passing through the sample and the local discharges between particles will melt superficially the particles external surface, but due to the very short time of pressing and passing the current through the powders mixture the melt will solidify rapidly and will assure a very strong bonding between the particles. In this way, large density nanocomposite magnets are produced.

Table 4 presents the magnetic characteristics of nanocrystalline FINEMET and amorphous CoFeMnSiB melt-spun ribbons as well as for the Nd-Fe-B/FINEMET and Nd-Fe-B/CoFeMnSiB SPS nanocomposite magnets. The magnetic properties of the SPS magnets are altered slightly compared with the Nd16 precursor powders due most probably by the mechanical deformations induced during SPS processing. Once the amount of FINEMET powders is increasing, the saturation magnetization increases because of the higher saturation value of FINEMET compared with Nd16 powders, whilst when increasing the CoFeMnSiB content the situation is opposite due to the lower saturation magnetization of Co-based amorphous alloy. However, the largest values for the coercive field and the maximum energy product are obtained for Nd16/CoFeMnSiB SPS magnets, as a consequence of the stronger exchange interactions between the nanograins of the soft and hard magnetic phases. For both types of SPS nanocomposite magnets the optimum magnetic properties are obtained for a sintering pressure of 60 MPa.

The coercive field is determined by two main factors: (1) the pinning of the domain walls and (2) the nucleation of the reversed magnetic domains. The simultaneous presence of the two mechanisms is confirmed by the shape of the first magnetization curve for Nd16/FINEMET SPS magnets prepared at two different sintering pressures and shown in Fig. 14. For the SPS magnets sintered at 50 MPa the coercive field is mainly determined by the nucleation mechanism, whereas for the larger sintering pressure of 60 MPa the pinning mechanism becomes predominant.
Table 4. Magnetic characteristics for nanocrystalline FINEMET and Nd16 ribbons, CoFeMnSiB amorphous ribbons as well as for SPS nanocomposite magnets prepared by combining the Nd16 hard magnetic and FINEMET or Co-based soft magnetic powders.

<table>
<thead>
<tr>
<th>Sample</th>
<th>( iH_c ) (kOe)</th>
<th>( 4\pi M_3T ) (kJG)</th>
<th>( 4\pi M_r ) (kJG)</th>
<th>( (BH)_{\text{max}} ) (MGOe)</th>
</tr>
</thead>
<tbody>
<tr>
<td>FINEMET ribbons</td>
<td>1.5x10^{-3}</td>
<td>12.2</td>
<td>9.5</td>
<td>--</td>
</tr>
<tr>
<td>CoFeSiBMn ribbons</td>
<td>5.0x10^{-3}</td>
<td>6.7</td>
<td>5.2</td>
<td>--</td>
</tr>
<tr>
<td>Nd\textsubscript{16}Fe\textsubscript{78}B\textsubscript{6} ribbons</td>
<td>20.0</td>
<td>8.9</td>
<td>6.9</td>
<td>11.9</td>
</tr>
<tr>
<td>NdFeB-4FINEMET / 50MPa</td>
<td>18.0</td>
<td>9.0</td>
<td>6.0</td>
<td>9.0</td>
</tr>
<tr>
<td>NdFeB-4FINEMET / 60MPa</td>
<td>18.5</td>
<td>9.0</td>
<td>6.3</td>
<td>9.9</td>
</tr>
<tr>
<td>NdFeB-5FINEMET / 50MPa</td>
<td>17.7</td>
<td>9.1</td>
<td>6.2</td>
<td>9.6</td>
</tr>
<tr>
<td>NdFeB-5FINEMET / 60MPa</td>
<td>18.0</td>
<td>9.1</td>
<td>6.3</td>
<td>9.9</td>
</tr>
<tr>
<td>NdFeB-6FINEMET / 50MPa</td>
<td>18.5</td>
<td>9.2</td>
<td>6.0</td>
<td>9.0</td>
</tr>
<tr>
<td>NdFeB-6FINEMET / 60MPa</td>
<td>18.0</td>
<td>9.2</td>
<td>5.9</td>
<td>8.7</td>
</tr>
<tr>
<td>NdFeB-4CoFeMnSiB / 50MPa</td>
<td>17.2</td>
<td>8.7</td>
<td>6.0</td>
<td>9.0</td>
</tr>
<tr>
<td>NdFeB-4CoFeMnSiB / 60MPa</td>
<td>19.5</td>
<td>8.4</td>
<td>6.1</td>
<td>9.3</td>
</tr>
<tr>
<td>NdFeB-5CoFeMnSiB / 50MPa</td>
<td>16.8</td>
<td>8.5</td>
<td>5.9</td>
<td>8.7</td>
</tr>
<tr>
<td>NdFeB-5CoFeMnSiB / 60MPa</td>
<td>18.2</td>
<td>8.8</td>
<td>6.4</td>
<td>10.2</td>
</tr>
<tr>
<td>NdFeB-6CoFeMnSiB / 50MPa</td>
<td>19.5</td>
<td>8.7</td>
<td>6.2</td>
<td>9.6</td>
</tr>
<tr>
<td>NdFeB-6CoFeMnSiB / 60MPa</td>
<td>17.7</td>
<td>8.7</td>
<td>6.3</td>
<td>9.9</td>
</tr>
</tbody>
</table>

Fig. 14. The first magnetization curves for Nd16-6\%FINEMET SPS magnets sintered at 2 different sintering pressures.

In Fig. 15 are presented comparatively the hysteresis loops for Nd16/FINEMET and Nd16/CoFeMnSiB SPS nanocomposite magnets sintered at the same sintering pressure (60 MPa). The remanence and coercive field for the 2 nanocomposite magnets are almost the same, but the presence of the shoulder in the second quadrant is almost not visible for the Nd16/CoFeMnSiB magnets because of the better exchange coupling of the soft and hard magnetic phases.

Fig. 16 shows the backscattered SEM micrographs of Nd16/5Fe SPS magnets consolidated at two pressures: 50 MPa (a) and 80 MPa (b). One can notice the difference in the morphologies of the two magnets and the fact that higher pressures are melting the powders.
boundaries and the resulting material is similar with a bulk one. Both magnets prepared at
50 and 80 MPa, respectively, are mainly consisting of primary Nd$_2$Fe$_{14}$B hard magnetic
phase and secondary αFe soft magnetic phase, as proved by XRD patterns shown in Fig. 8.

Fig. 16. BSE-SEM micrographs of Nd16Fe78B6+5 wt.% Fe SPS magnets consolidated at: (a)
50 MPa and (b) 80 MPa.

The magnetic characteristics of Nd-Fe-B/Fe SPS nanocomposite magnets are shown in Table
5. The remanence increases from 6.4 kG for Nd15 SPS magnets without addition of Fe to 6.6
kG for Nd15/5 wt.% Fe SPS magnets consolidated at the same pressure of 50 MPa. Further
increase of Fe powders content leads to the decrease of the remanence to 5.6 kG for Nd15/10
wt.% Fe SPS magnets consolidated at 50 MPa. The enhancement of the remanence of the SPS
nanocomposite magnets comes from the exchange interactions between the powders of the
hard and soft magnetic phases, respectively. The increase of the Fe powders content leads to
a slight decrease of the coercivity, too. A more evident proof of the increasing role of
exchange interactions in SPS nanocomposite magnets appears when comparing the melt-
spun ribbons, the precursor powders and the SPS samples. Whereas in the ribbons the
coupling between the grains of the hard and soft magnetic phases is very good, it is
deteriorating for the powders and the decoupling shoulder appears on the M-H loops. The
decoupling is becoming less evident for SPS nanocomposite magnets with 4 and 5 wt.% Fe,
and is reappearing for larger contents of Fe powders. The coercive field also increases with
the Fe powders content to 5 wt.%, reaching almost the value obtained for melt-spun ribbons.
in the as-quenched state. A similar behaviour is observed for Nd16 SPS consolidated magnets. The magnetic properties are altered for all SPS nanocomposite magnets consolidated at higher pressures (80 MPa). Correlating the microstructure evolution with the observed changes in their magnetic properties, the alteration could be determined by the effect of both the distance modification between particles and their deformation (see Fig. 16, too). By consolidation through the SPS technique the maximum energy product generally increases for optimal processing parameters, the highest increase of about 8 % being obtained for Nd15/5 wt.% Fe SPS nanocomposite magnets consolidated at 50 MPa compared with the similar Nd15 SPS magnets prepared in the same conditions.

<table>
<thead>
<tr>
<th>Sample</th>
<th>$H_c$ (kOe)</th>
<th>$4\pi M_r$ (kG)</th>
<th>$(BH)_{max}$ (MGOe)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Nd15Fe9B6</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Nd15 / 50MPa</td>
<td>17.8</td>
<td>6.4</td>
<td>10.2</td>
</tr>
<tr>
<td>Nd15-4Fe / 50MPa</td>
<td>17.4</td>
<td>6.6</td>
<td>10.9</td>
</tr>
<tr>
<td>Nd15-4Fe / 80MPa</td>
<td>17.2</td>
<td>6.4</td>
<td>10.2</td>
</tr>
<tr>
<td>Nd15-5Fe / 50MPa</td>
<td>18.1</td>
<td>6.6</td>
<td>10.9</td>
</tr>
<tr>
<td>Nd15-5Fe / 80MPa</td>
<td>17.2</td>
<td>6.5</td>
<td>10.6</td>
</tr>
<tr>
<td>Nd15-10Fe / 50MPa</td>
<td>16.8</td>
<td>6.0</td>
<td>9.0</td>
</tr>
<tr>
<td>Nd15-10Fe / 80MPa</td>
<td>15.8</td>
<td>5.6</td>
<td>7.8</td>
</tr>
<tr>
<td>Nd16Fe9B6</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Nd16 / 50MPa</td>
<td>17.5</td>
<td>6.1</td>
<td>9.3</td>
</tr>
<tr>
<td>Nd16-4Fe / 50MPa</td>
<td>17.4</td>
<td>6.4</td>
<td>10.2</td>
</tr>
<tr>
<td>Nd16-4Fe / 80MPa</td>
<td>16.7</td>
<td>6.2</td>
<td>9.6</td>
</tr>
<tr>
<td>Nd16-5Fe / 50MPa</td>
<td>16.7</td>
<td>6.4</td>
<td>10.2</td>
</tr>
<tr>
<td>Nd16-5Fe / 80MPa</td>
<td>16.5</td>
<td>5.9</td>
<td>8.7</td>
</tr>
<tr>
<td>Nd16-10Fe / 50MPa</td>
<td>15.5</td>
<td>6.2</td>
<td>9.6</td>
</tr>
<tr>
<td>Nd16-10Fe / 80MPa</td>
<td>15.0</td>
<td>5.9</td>
<td>8.7</td>
</tr>
</tbody>
</table>

Table 5. Magnetic characteristics for Nd15, Nd16, Nd15/Fe and Nd16/Fe SPS nanocomposite magnets.

Fig. 17 presents comparatively the $\delta M$ curves for Nd15/Fe SPS nanocomposite magnets consolidated at 50 MPa. When the applied field is small and the nanocomposite magnet is in the initial magnetization state the magnetic moments within the nanograins are randomly oriented and the distribution of angles between adjacent moments is broad. Hence, there are exchange interactions between nanograins and $\delta M$ has positive values (the ascending part of the $\delta M$ plots). The increase of the applied field favors the alignment of the neighboring magnetic moments in the nanograins and the exchange interactions become predominant. Consequently, $\delta M$ increases and reaches a peak. The shift of the maximum field value to larger fields for 4 wt.% Fe and 5 wt.% SPS nanocomposite magnets is also consistent with the increase in coercivity (Table 5). Additionally, a higher $\delta M$ peak indicates stronger exchange interactions between the nanograins of the magnetically hard and soft phases. The strongest exchange interactions are observed for a content of 5 wt. % Fe. While the applied field approaches the coercivity (Fig. 17 and Table 5), the moments basically cancel out each other and the effective magnetization in the direction of applied field approximates to zero. Thus, the magnetostatic interactions are still very weak and the exchange interactions are
predominant giving positive $\delta M$ values. With the continuous increase of the applied field the magnetic moments are gradually aligned parallel in each grain and the dipolar interactions between nanograins are becoming more important compared with the exchange interactions. Consequently, $\delta M$ becomes slightly negative.

Fig. 17. $\delta M$ plots for Nd15/Fe SPS nanocomposite magnets consolidated at 50 MPa.

M-H loops for Nd16/FeCo SPS nanocomposite magnets presented in Fig. 18 show that the largest coercive field is reached for 96Nd16:4FeCo. For this ratio the exchange coupling is the strongest and no decoupling is visible in the second quadrant. By increasing the FeCo content the decoupling between soft and hard magnetic nanograins is evident and the shoulder appearing on the demagnetization curves is more pronounced.

Fig. 18. Magnetic hysteresis loops for Nd16/FeCo SPS nanocomposite magnets sintered at 50 MPa. The hysteresis loop for the precursor Nd16 powders is shown for comparison.

Whereas the coercive field values are increasing with the increase of Nd content in the hard magnetic phase precursors, the largest values of the maximum energy product are achieved for Nd12/4FeCo SPS nanocomposite magnets (Table 6).
Table 6. The magnetic characteristics of Nd-Fe-B/FeCo SPS nanocomposite magnets sintered at 50 MPa.

<table>
<thead>
<tr>
<th>Sample</th>
<th>$H_c$ (kOe)</th>
<th>$4\pi M_s$ (kG)</th>
<th>$4\pi M_r$ (kG)</th>
<th>$(BH)_{max}$ (MGOe)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Nd8-4FeCo</td>
<td>4.0</td>
<td>10.8</td>
<td>5.7</td>
<td>8.1</td>
</tr>
<tr>
<td>Nd8-10FeCo</td>
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<td>13.3</td>
<td>7.1</td>
<td>12.6</td>
</tr>
<tr>
<td>Nd10-4FeCo</td>
<td>5.2</td>
<td>10.1</td>
<td>6.2</td>
<td>9.6</td>
</tr>
<tr>
<td>Nd10-10FeCo</td>
<td>4.7</td>
<td>12.8</td>
<td>7.2</td>
<td>12.9</td>
</tr>
<tr>
<td>Nd12-4FeCo</td>
<td>10.0</td>
<td>10.4</td>
<td>7.6</td>
<td>14.4</td>
</tr>
<tr>
<td>Nd12-10FeCo</td>
<td>8.4</td>
<td>11.5</td>
<td>6.6</td>
<td>10.9</td>
</tr>
<tr>
<td>Nd15-4FeCo</td>
<td>13.0</td>
<td>9.1</td>
<td>6.5</td>
<td>10.5</td>
</tr>
<tr>
<td>Nd15-10FeCo</td>
<td>12.0</td>
<td>9.6</td>
<td>5.5</td>
<td>7.6</td>
</tr>
<tr>
<td>Nd16-4FeCo</td>
<td>17.6</td>
<td>8.4</td>
<td>5.8</td>
<td>8.4</td>
</tr>
<tr>
<td>Nd16-10FeCo</td>
<td>13.9</td>
<td>9.0</td>
<td>5.7</td>
<td>8.1</td>
</tr>
</tbody>
</table>

Fig. 19. $\delta M(H)$ curves as a function of Nd content in the hard magnetic component precursors. All SPS nanocomposites magnets contain the same amount of Fe\textsubscript{2}Co powders, i.e. 4 wt.%, and they are consolidated at 50 MPa.

For nanocomposite magnets with low contents of Nd (Nd8/4FeCo and Nd10/4FeCo) the values of $\delta M$ shown in Fig. 19 are negative suggesting that the interactions between the soft and hard magnetic nanograins have dipolar (long-range magnetostatic) character and promote the demagnetized state. The magnetostatic interactions reduce the number of irreversible magnetization processes during magnetization in comparison with those occurring during demagnetization. The positive $\delta M$ curves for Nd-rich (Nd15/4FeCo and Nd16-4FeCo) nanocomposite magnets suggest stronger exchange interactions dominating the magnetization processes. $\delta M$ increases for Nd12/4FeCo nanocomposite magnets and exhibits a peak on the positive side when increasing the applied field, then decreases to zero and changes to negative for much larger applied fields than the coercive field (Fig. 19 and Table 6).

The highest positive peak is obtained for Nd15/4FeCo nanocomposite magnets, suggesting predominant exchange interactions in the detriment of the magnetostatic ones. The same
behaviour is observed independent of the FeCo content in the Nd-Fe-B/FeCo SPS nanocomposite magnets.

Fig. 20 presents the first magnetization curves for Nd8/4FeCo, Nd10/4FeCo and Nd12/4FeCo SPS nanocomposite magnets sintered at 50 MPa.

![First magnetization curves for NdFeB/FeCo SPS nanocomposite magnets consolidated at 50 MPa.](image)

From the shape of the first magnetization curves we can conclude that for Nd8/4FeCo and Nd10/4FeCo magnets the nucleation of the domain walls with reversed magnetization is predominant, whereas the mechanism is changed to the domain wall pinning for Nd12/4FeCo. This is related also to the fact that in Nd12/4FeCo nanocomposite magnets the amount of 2:14:1 phase is larger than for the sub-stoichiometric (Nd 8 and Nd10) compositions or Nd-rich compositions (Nd 15 and Nd16). The increase of the sintering temperature from 600°C to 650°C or of the sintering time from 5 to 7 min. results in the deterioration of the magnetic properties due to the increased amount of soft magnetic phases as well as to the increase of the nanograins size from 20-30 nm to 150-200 nm. A temperature below 600°C or a sintering time shorter than 5 min. are not enough for a good bonding between the hard and soft magnetic components, even at a higher sintering pressure of 80 MPa. The increase of the sintering pressure from 50 to 80 MPa increases a little bit the remanence, but leaves the coercive field almost constant.

### 6. Conclusions and future perspectives

Nanocomposite permanent magnets have been obtained by SPS consolidation of powders of different hard and soft magnetic materials in different ratios. The powders have been obtained by ball-milling the NdFeB hard magnetic nanocrystalline melt-spun ribbons precursors and respectively FeSiB or CoFeSiB amorphous wires, FINEMET nanocrystalline melt-spun ribbons, CoFeMnSiB amorphous melt-spun ribbons, Fe2Co fully crystalline melt-spun ribbons or Fe commercial micropowders.

Despite the fact that the powders exhibit larger coercive fields, the coupling between the soft and hard magnetic phases is improved for nanocomposite magnets. The strength of the decoupling is indicated by the second shoulder of dM/dH=f(H) curves located nearby Y
axis, which is very small for SPS consolidated nanocomposite magnets. Most probably a collective process of both exchange and magnetostatic interactions between the powders of the 2 materials forming the nanocomposite is responsible for such specific behaviour. The exchange coupling plays the predominant role in Nd-Fe-B melt-spun nanocrystalline ribbons, whereas the magnetostatic interactions becomes more predominant for Nd-rich nanocomposites. However, the percentage of exchange interactions contribution to the macroscopic response of the nanocomposite systems is still very important, considering that usually the magnetostatic interactions are reducing the coercive field values.

The magnetic properties can be tailored by choosing the proper SPS consolidation conditions and the optimum ratios of the 2 or more different phases forming the nanocomposites.

Further studies on increasing the amount of soft magnetic component and the sintering pressure are necessary to increase the maximum energy product and to reach at least the values obtained for existing sintered magnets (maximum 55 MGOe), if not the calculated ones (64 MGOe) (Sagawa et al., 1985). To achieve this goal, the preparation of anisotropic SPS nanocomposite magnets consisting of Nd-Fe-B hard magnetic and Fe or Co-Fe based soft magnetic materials is imperiously necessary. Another goal is to increase the operating temperature of these magnets from 100-150°C to at least 250-300°C by adding to the Nd-Fe-B hard magnetic component Co-based materials or other additives which could increase the Curie temperature of these SPS nanocomposite PMs.

7. Acknowledgments

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8. References


Advances in Nanocomposites - Synthesis, Characterization and Industrial Applications

Edited by Dr. Boreddy Reddy

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Advances in Nanocomposites - Synthesis, Characterization and Industrial Applications was conceived as a comprehensive reference volume on various aspects of functional nanocomposites for engineering technologies. The term functional nanocomposites signifies a wide area of polymer/material science and engineering, involving the design, synthesis and study of nanocomposites of increasing structural sophistication and complexity useful for a wide range of chemical, physicochemical and biological/biomedical processes. "Emerging technologies" are also broadly understood to include new technological developments, beginning at the forefront of conventional industrial practices and extending into anticipated and speculative industries of the future. The scope of the present book on nanocomposites and applications extends far beyond emerging technologies. This book presents 40 chapters organized in four parts systematically providing a wealth of new ideas in design, synthesis and study of sophisticated nanocomposite structures.

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